

Ensemble density functional theory of electrons and nuclei

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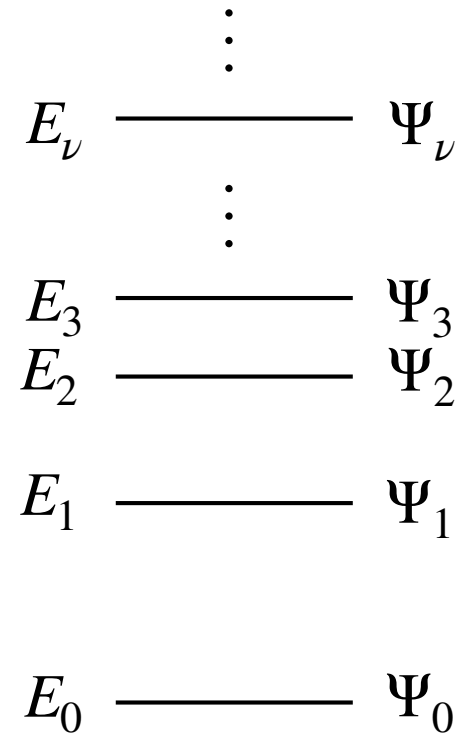
*Brief review of (ensemble) **density-functional**
electronic structure **theory***

Density functional theory (DFT) of ground and excited states

$$\begin{array}{ccc} & \vdots & \\ E_\nu & \text{-----} & \Psi_\nu \\ & \vdots & \\ E_3 & \text{-----} & \Psi_3 \\ E_2 & \text{-----} & \Psi_2 \\ \\ E_1 & \text{-----} & \Psi_1 \\ \\ E_0 & \text{-----} & \Psi_0 \end{array}$$

Density functional theory (DFT) of ground and excited states

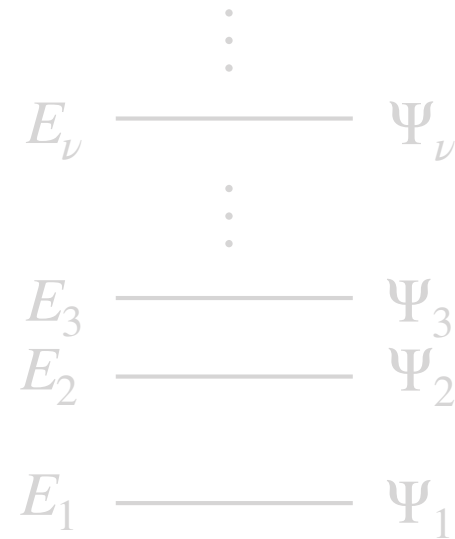
$$\hat{H}\Psi_\nu = E_\nu\Psi_\nu$$



Density-functional theory (DFT) of N -electron ground states

$$E_0 = \min_{\Psi} \langle \Psi | \hat{H} | \Psi \rangle$$

Variational principle

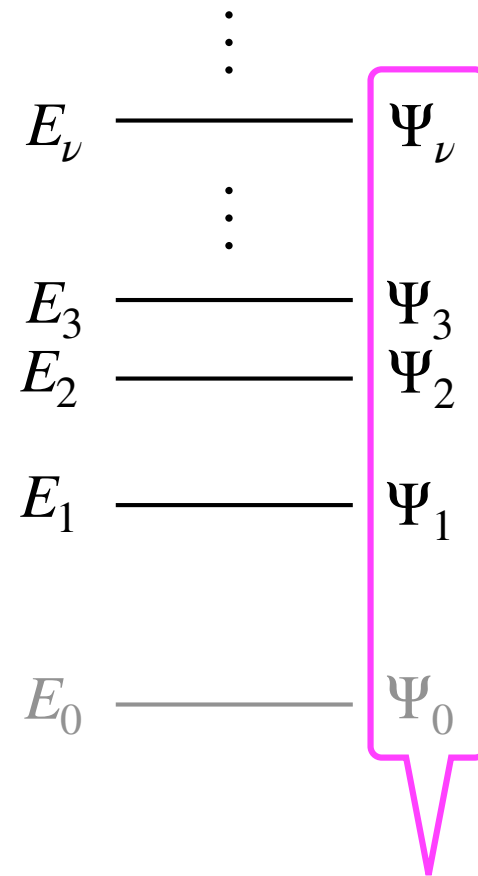


$$E_0 \text{ ————— } \Psi_0$$

$$n_{\Psi_0}(\mathbf{r}) \equiv n_0^N(\mathbf{r})$$

Ground-state
electronic **density**

Theophilou-Gross-Oliveira-Kohn (TGOK) ensembles



N -electron states

A. K. Theophilou, *J. Phys. C: Solid State Phys.* **12**, 5419 (1979).

A. K. Theophilou, in *The Single Particle Density in Physics and Chemistry*, edited by N. H. March and B. M. Deb (Academic Press, 1987), pp. 210–212.

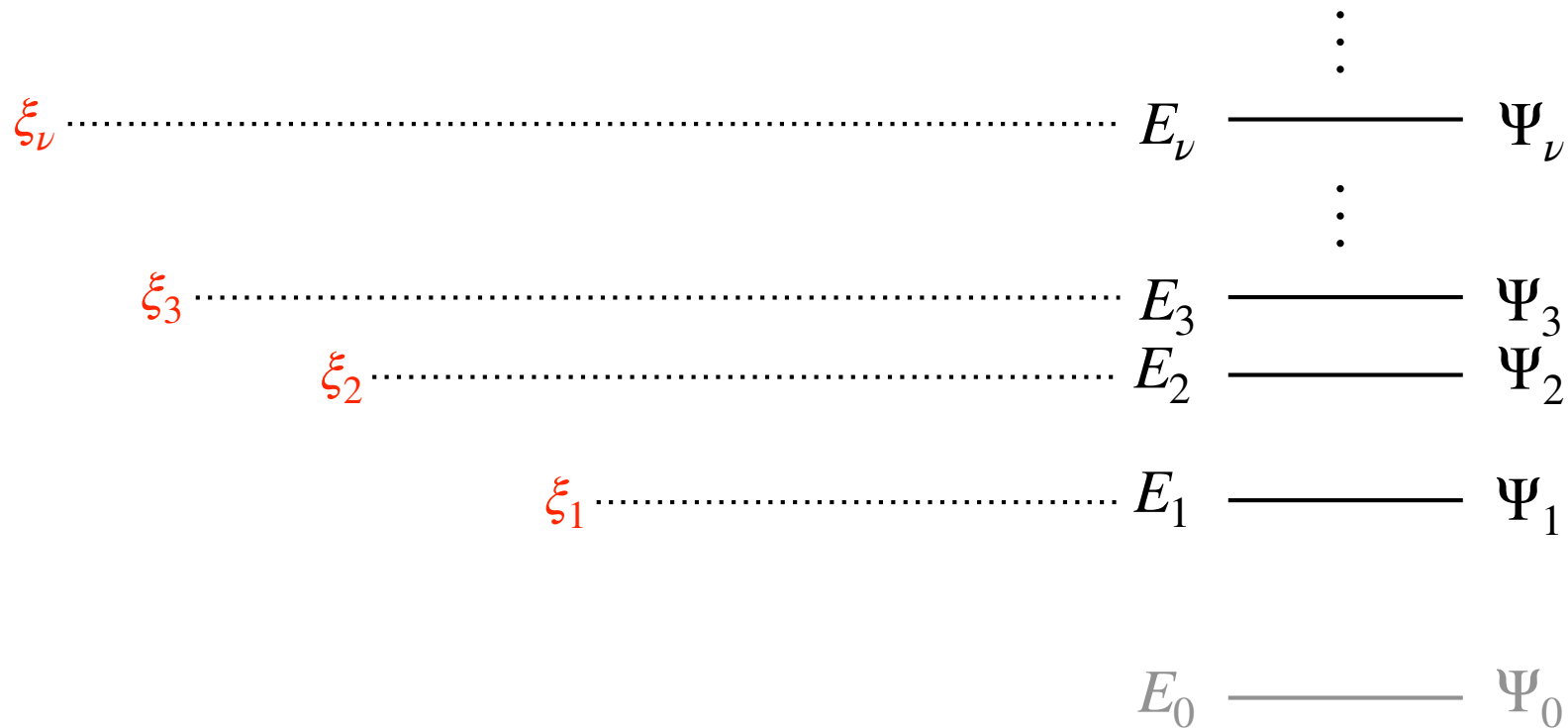
E. K. U. Gross, L. N. Oliveira, and W. Kohn, *Phys. Rev. A* **37**, 2805 (1988).

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L. N. Oliveira, E. K. U. Gross, and W. Kohn, *Phys. Rev. A* **37**, 2821 (1988).

K. Deur and E. Fromager, *J. Chem. Phys.* **150**, 094106 (2019).

Theophilou-Gross-Oliveira-Kohn (TGOK) ensembles



(Almost) independent **ensemble weights** $\{\xi_\nu\}$
are assigned to the excited states

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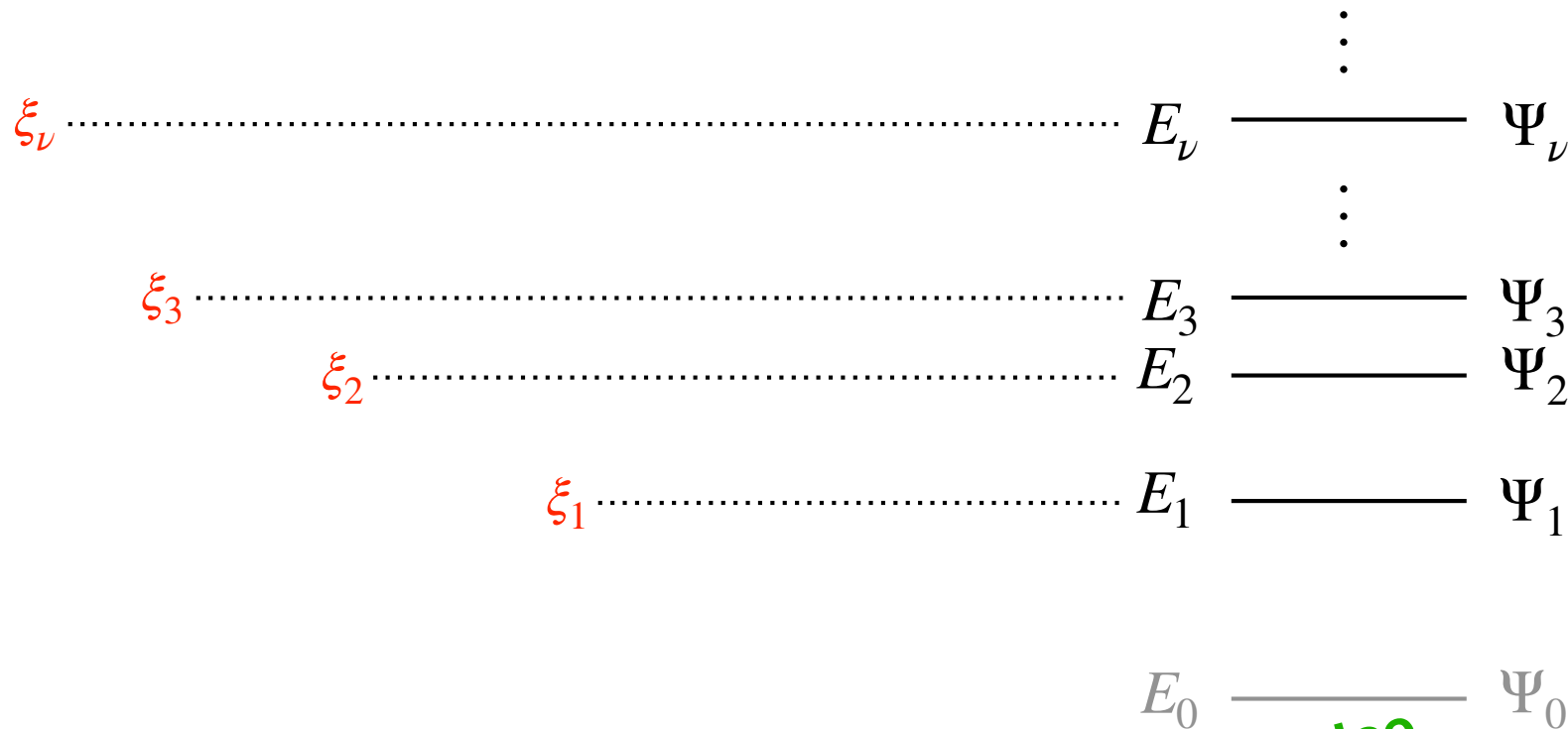
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Theophilou-Gross-Oliveira-Kohn (TGOK) ensembles



(Almost) independent **ensemble weights** $\{\xi_\nu\}$
are assigned to the excited states

**Arbitrary choice
in practice!**

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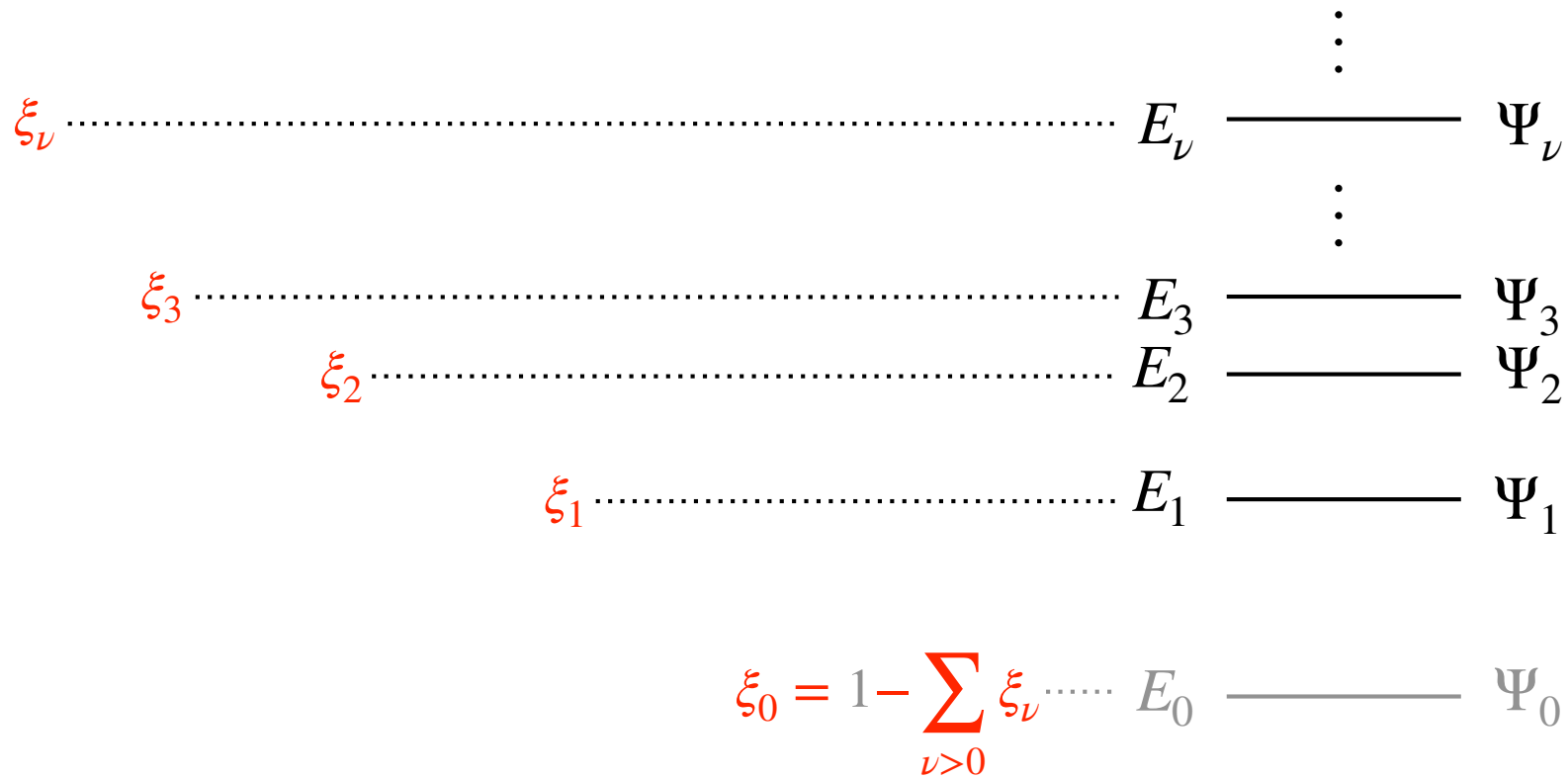
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Theophilou-Gross-Oliveira-Kohn (TGOK) ensembles



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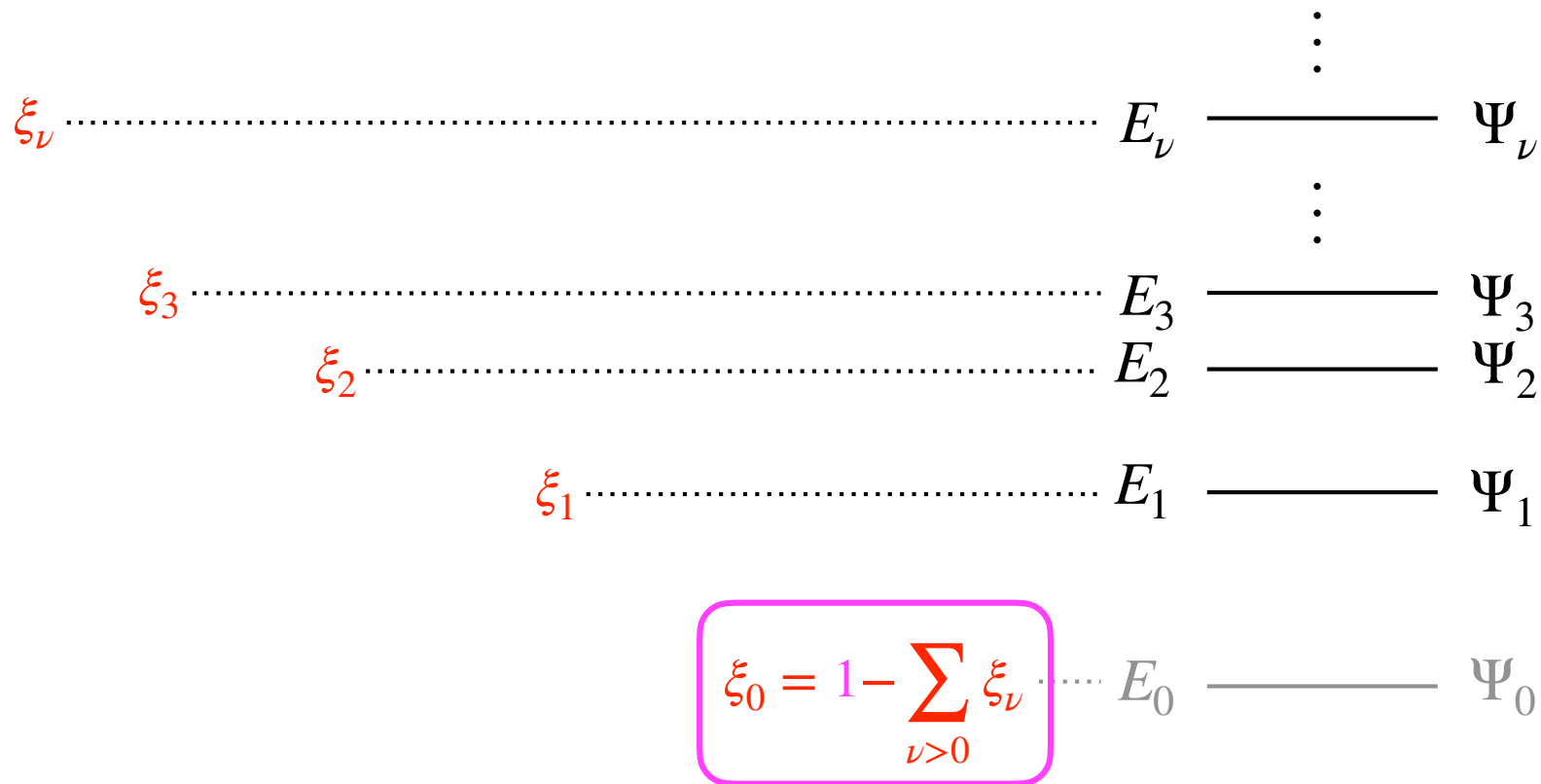
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Theophilou-Gross-Oliveira-Kohn (TGOK) ensembles



Neutral excitation processes

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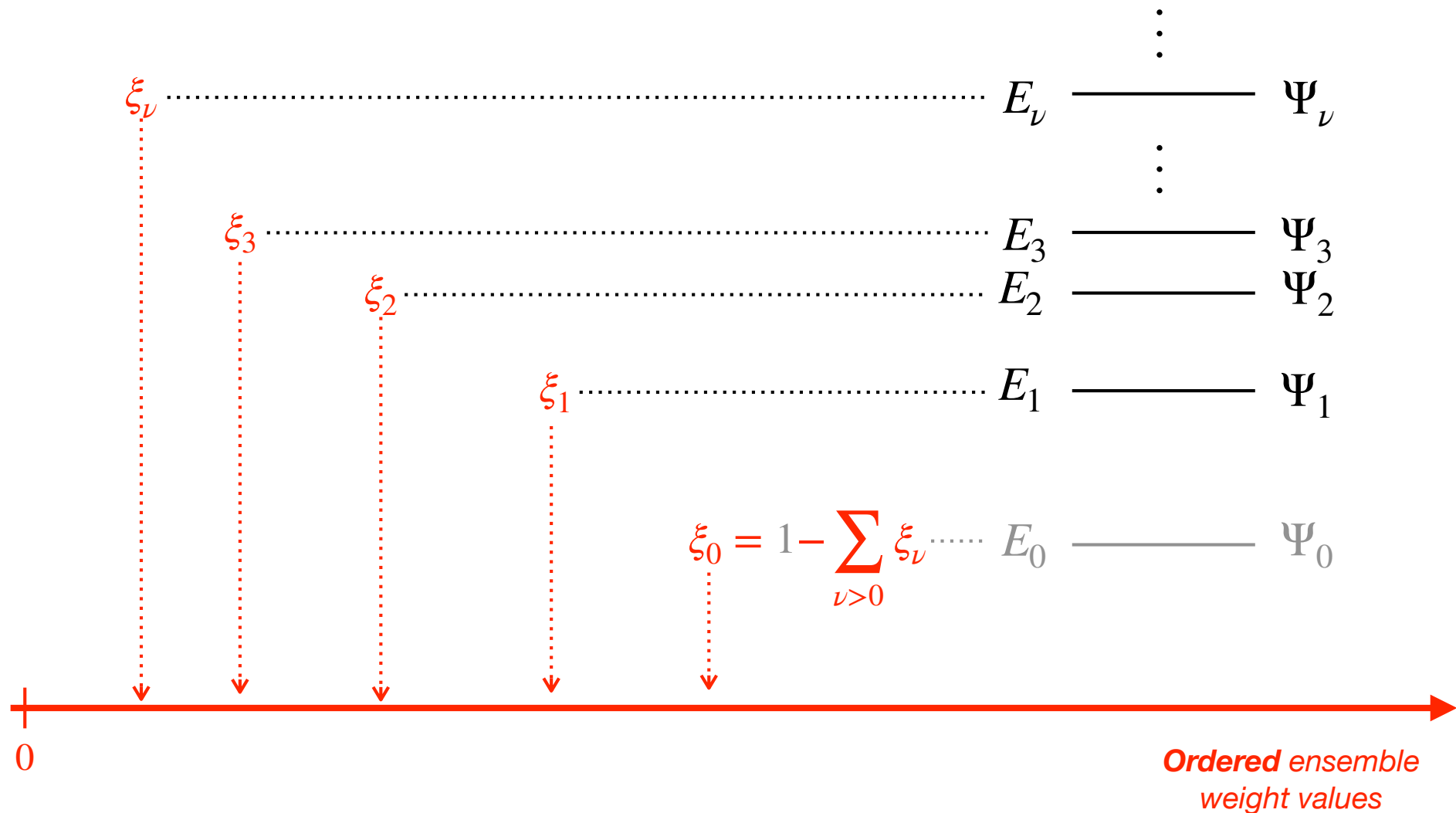
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Theophilou-Gross-Oliveira-Kohn (TGOK) ensembles



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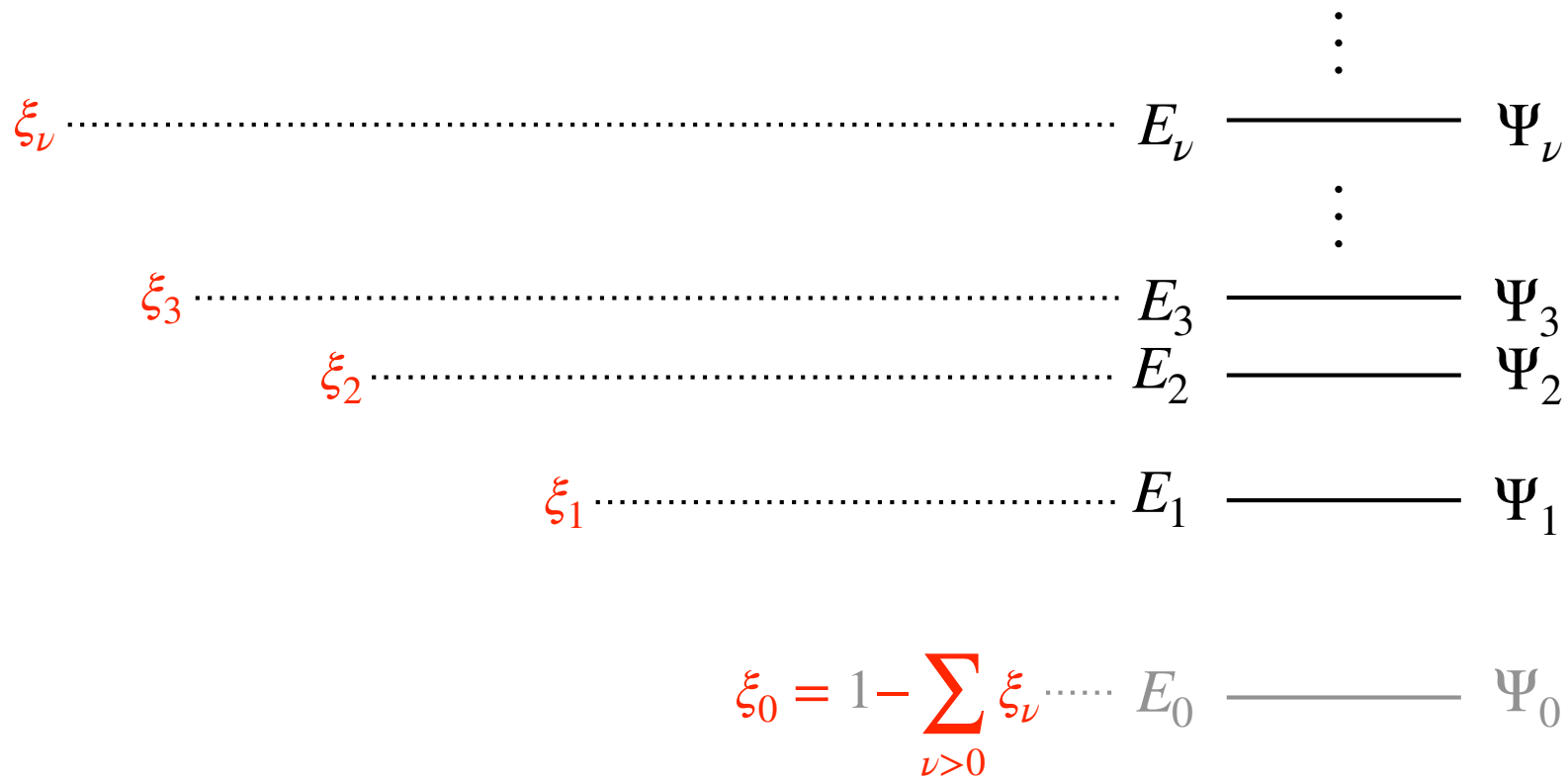
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Theophilou-Gross-Oliveira-Kohn (TGOK) variational principle

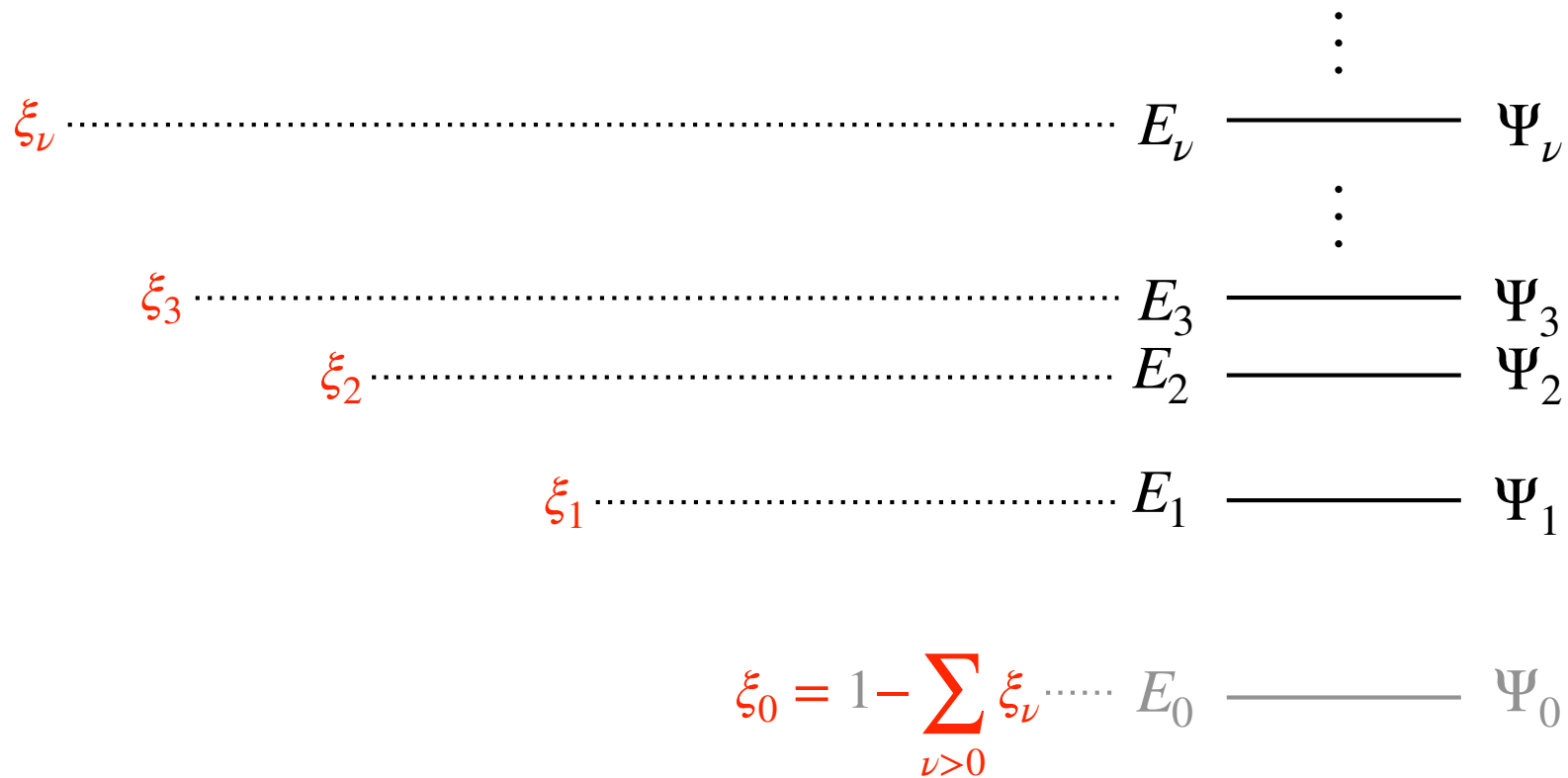


$$\sum_{\nu \geq 0} \xi_\nu \langle \tilde{\Psi}_\nu | \hat{H} | \tilde{\Psi}_\nu \rangle \geq \sum_{\nu \geq 0} \xi_\nu E_\nu$$

$\langle \tilde{\Psi}_\mu | \tilde{\Psi}_\nu \rangle = \delta_{\mu\nu}$

Variational principle for ensembles

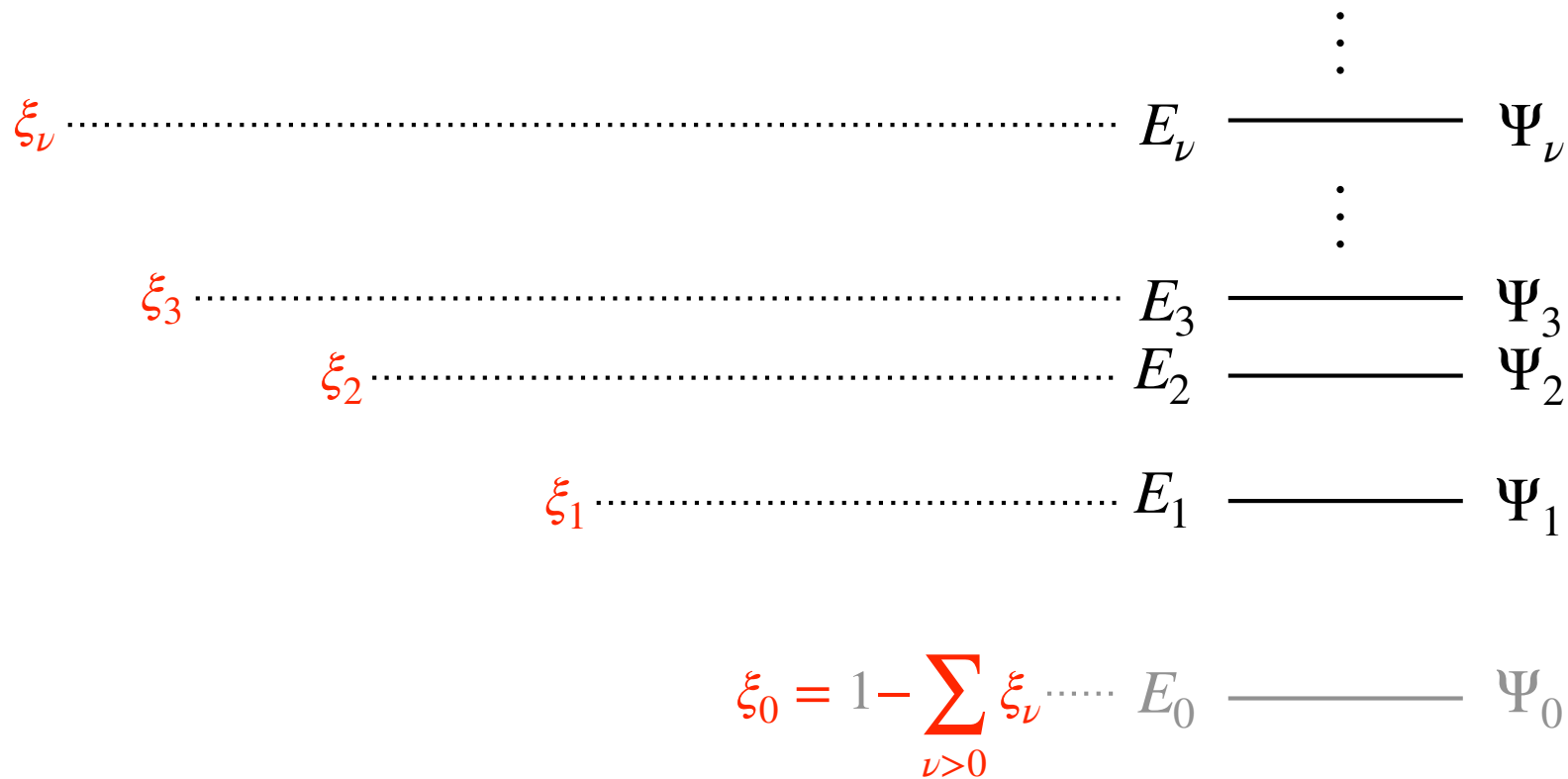
Theophilou-Gross-Oliveira-Kohn (TGOK) variational principle



$$\sum_{\nu \geq 0} \xi_\nu \langle \tilde{\Psi}_\nu | \hat{H} | \tilde{\Psi}_\nu \rangle \geq \sum_{\nu \geq 0} \xi_\nu E_\nu \equiv E\{\xi_\nu\}$$

Ensemble energy

Theophilou-Gross-Oliveira-Kohn (TGOK) variational principle



$$\sum_{\nu \geq 0} \xi_\nu \langle \tilde{\Psi}_\nu | \hat{H} | \tilde{\Psi}_\nu \rangle \geq \sum_{\nu \geq 0} \xi_\nu E_\nu \equiv E\{\xi_\nu\}[n] \leftarrow n(\mathbf{r}) \equiv \sum_{\nu \geq 0} \xi_\nu n_{\Psi_\nu}(\mathbf{r})$$

Ensemble energy
Ensemble density

How different is an **ensemble Kohn–Sham DFT** calculation from a regular (ground-state) one?

How different is an **ensemble Kohn–Sham DFT** calculation from a regular (ground-state) one?

$$n(\mathbf{r}) \equiv \sum_i \left(\sum_{\nu \geq 0} n_{i,\nu} \xi_\nu \right) \left| \varphi_i^{\{\xi_\nu\}}(\mathbf{r}) \right|^2$$

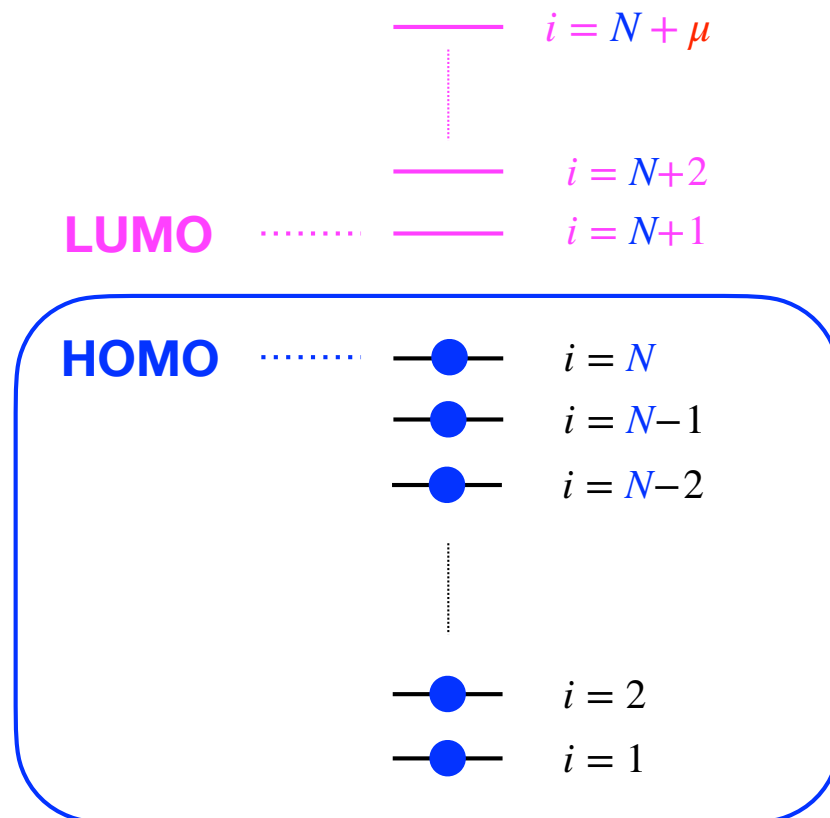
Fractionally occupied KS orbitals

How different is an **ensemble Kohn–Sham DFT** calculation from a regular (ground-state) one?

$$E_{\text{Hxc}}[n] \longrightarrow E_{\text{Hxc}}^{\{\xi_\nu\}}[n]$$

*Ensemble-weight
dependence*

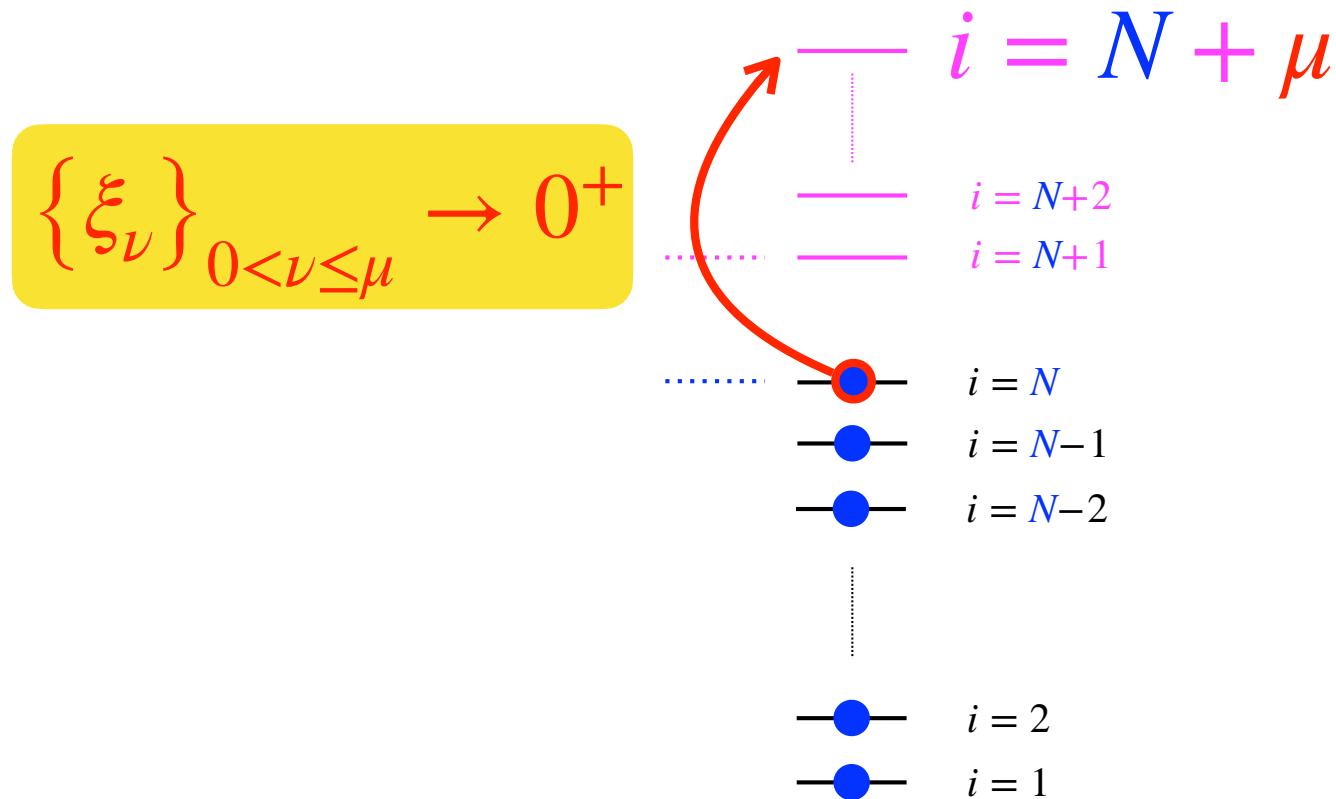
Ensemble weight dependence of the xc functional



$$\left\{ \begin{array}{l} \varepsilon \\ \xi_\nu \end{array} \right\}_{0 < \nu \leq \mu} = 0$$

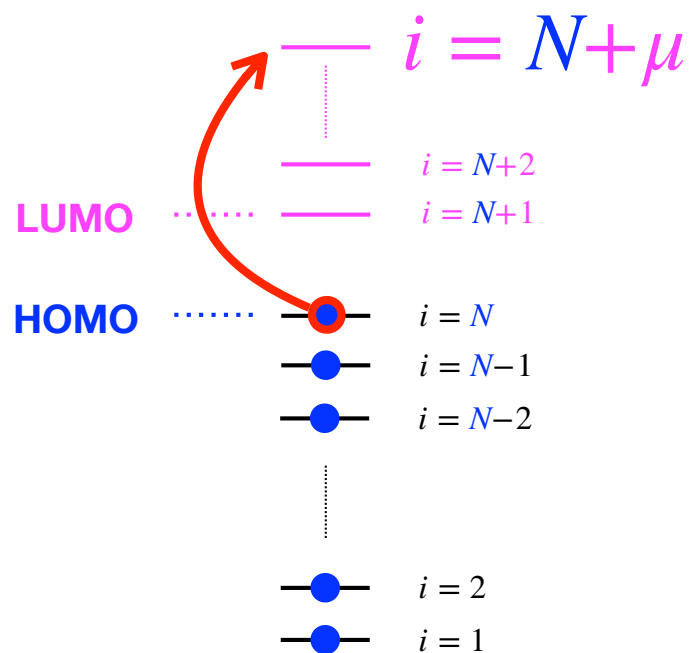
Regular **ground-state** Kohn–Sham configuration

Ensemble weight dependence of the xc functional



Infinitesimal excitation process in the fictitious Kohn–Sham world

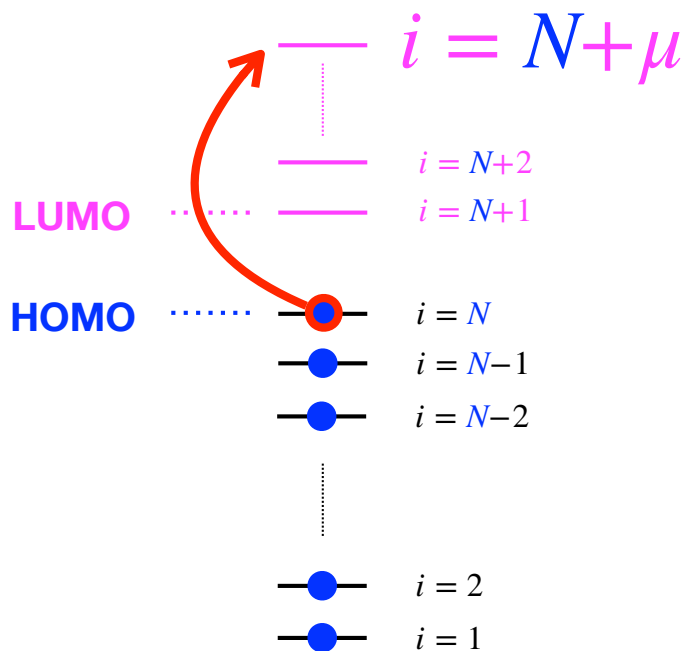
Ensemble weight dependence of the xc functional



$$E_{\mu} - E_0 = \varepsilon_{N+\mu}^{\{\xi_{\nu}\} \rightarrow 0^+} - \varepsilon_N^{\{\xi_{\nu}\} = 0}$$

Exact excitation energy!

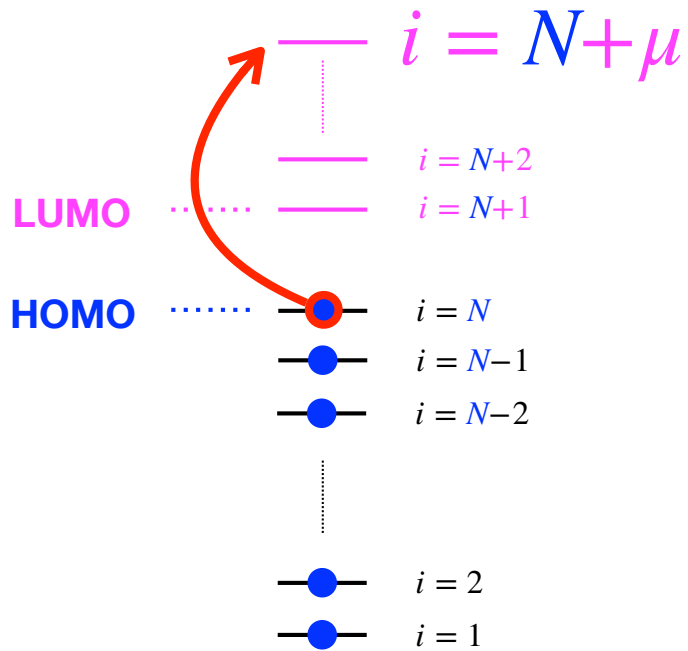
Ensemble weight dependence of the xc functional



$$E_{\mu} - E_0 = \varepsilon_{N+\mu}^{\{\xi_{\nu}\} \rightarrow 0^+} - \varepsilon_N^{\{\xi_{\nu}\} = 0}$$

$$v_{\text{Hxc}}^{\{\xi_{\nu}\}}(\mathbf{r}) = \frac{\delta E_{\text{Hxc}}^{\{\xi_{\nu}\}}[n]}{\delta n(\mathbf{r})}$$

Ensemble weight dependence of the xc functional



$$E_{\mu} - E_0 = \varepsilon_{N+\mu}^{\{\xi_{\nu}\} \rightarrow 0^+} - \varepsilon_N^{\{\xi_{\nu}\} = 0}$$

$$v_{\text{Hxc}}^{\{\xi_{\nu}\}}(\mathbf{r}) = \frac{\delta E_{\text{Hxc}}^{\{\xi_{\nu}\}}[n]}{\delta n(\mathbf{r})}$$

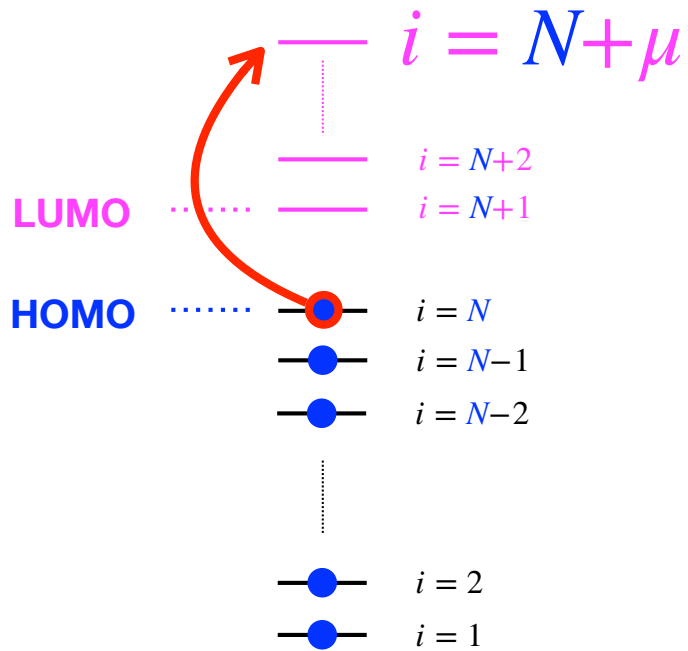
$$\int \frac{d\mathbf{r}}{N} \left(v_{\text{Hxc}}^{\{\xi_{\nu}\} \rightarrow 0^+}(\mathbf{r}) - v_{\text{Hxc}}^{\{\xi_{\nu}\} = 0}(\mathbf{r}) \right) n(\mathbf{r}) \Big|_{n=n_0^N} = ???$$

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Ensemble weight dependence of the xc functional



$$E_{\mu} - E_0 = \varepsilon_{N+\mu}^{\{\xi_{\nu}\} \rightarrow 0^+} - \varepsilon_N^{\{\xi_{\nu}\} = 0}$$

$$v_{\text{Hxc}}^{\{\xi_{\nu}\}}(\mathbf{r}) = \frac{\delta E_{\text{Hxc}}^{\{\xi_{\nu}\}}[n]}{\delta n(\mathbf{r})}$$

$$\int \frac{d\mathbf{r}}{N} \left(v_{\text{Hxc}}^{\{\xi_{\nu}\} \rightarrow 0^+}(\mathbf{r}) - v_{\text{Hxc}}^{\{\xi_{\nu}\} = 0}(\mathbf{r}) \right) n(\mathbf{r}) \Big|_{n=n_0^N} = \frac{\partial E_{\text{Hxc}}^{\{\xi_{\nu}\}}[n]}{\partial \xi_{\mu}} \Big|_{\{\xi_{\nu}\} = 0}$$

“Derivative discontinuity”

*Towards an **in-principle-exact** ensemble
density-functional theory of **electrons and nuclei***

Exact Factorization-Based Density Functional Theory of Electrons and Nuclei

Ryan Requist* and E. K. U. Gross

Max Planck Institute of Microstructure Physics, Weinberg 2, 06120 Halle, Germany

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Direct Nonadiabatic Dynamics by Mixed Quantum-Classical Formalism Connected with Ensemble Density Functional Theory Method: Application to *trans*-Penta-2,4-dieniminium Cation

Michael Filatov,*^{ID} Seung Kyu Min,*^{ID} and Kwang S. Kim^{ID}

Department of Chemistry, School of Natural Sciences, Ulsan National Institute of Science and Technology (UNIST), Ulsan 44919, Korea

Born–Huang expansion of the molecular wave function

$$\psi^{mol.}(R, r) = \sum_{\nu \geq 0} \chi_{\nu}(R) \Psi_{\nu}^{elec.}(R, r)$$

Trial molecular wave function

Bond distance (for example)

Electronic coordinates

Nuclear wave functions

Ground- and excited-state electronic wave functions

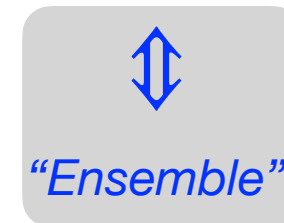
Born–Huang expansion of the molecular wave function

$$\psi^{mol.}(R, r) = \sum_{\nu \geq 0} \chi_{\nu}(R) \Psi_{\nu}^{elec.}(R, r)$$

*Trial molecular
wave function*

*Nuclear
wave functions*

*Ground- and excited-state
electronic wave functions*



Molecular variational principle

diatomics

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) |\underline{\chi}(R)|^2 \right. \\ \left. -\frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \underline{\Lambda}_{\underline{\Psi}^{elec.}}^{(j)}(R) \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \int dR |\underline{\chi}(R)|^2 \sum_{\nu \geq 0} \xi_\nu(R) E_\nu^{elec.}(R) \right\}$$

$$\psi^{mol.}(R, r) = \underline{\chi}^T(R) \underline{\Psi}^{elec.}(R, r)$$

The electronic energy within the molecule is an **ensemble energy**

Molecular variational principle

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) |\underline{\chi}(R)|^2 \right. \\ \left. -\frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \underline{\Lambda}_{\underline{\Psi}^{elec.}}^{(j)}(R) \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \int dR |\underline{\chi}(R)|^2 \sum_{\nu \geq 0} \xi_\nu(R) E_\nu^{elec.}(R) \right\}$$

To-be-determined ensemble weights:

$$\xi_\nu(R) \equiv \frac{|\chi_\nu(R)|^2}{|\underline{\chi}(R)|^2}$$



Molecular variational principle

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) \left| \underline{\chi}(R) \right|^2 \right. \\ \left. - \frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \Lambda_{\underline{\Psi}^{elec.}}^{(j)}(R) \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \dots \right.$$

$$\left\{ \int dr \Psi_{\nu}^{elec.}(R, r) \frac{\partial^j \Psi_{\mu}^{elec.}(R, r)}{\partial R^j} \right\} \equiv \Lambda_{\underline{\Psi}^{elec.}}^{(j)}(R)$$

Non-adiabatic couplings (NACs)

Molecular variational principle

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) \left| \underline{\chi}(R) \right|^2 \right. \\ \left. -\frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \underbrace{\Lambda_{\underline{\Psi}^{elec.}}^{(j)}(R)}_{\equiv \underline{\Psi}^{elec.}} \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \dots \right\}$$

$$\left\{ \int dr \underline{\Psi}_{\nu}^{elec.}(R, r) \frac{\partial^j \underline{\Psi}_{\mu}^{elec.}(R, r)}{\partial R^j} \right\} \equiv \underline{\Lambda}_{\underline{\Psi}^{elec.}}^{(j)}(R) = ???$$



Non-adiabatic couplings (NACs)

Molecular variational principle

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) \left| \underline{\chi}(R) \right|^2 \right. \\ \left. - \frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \Lambda_{\underline{\Psi}^{elec.}}^{(j)}(R) \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \dots \right\}$$

$$\left\{ \int dr \Psi_{\nu}^{elec.}(R, r) \frac{\partial^j \Psi_{\mu}^{elec.}(R, r)}{\partial R^j} \right\} \equiv \Lambda_{\underline{\Psi}^{elec.}}^{(j)}(R) = ??? \neq \Lambda_{\underline{\Phi}^{KS}}^{(j)}(R)$$

True non-adiabatic couplings (NACs)

Kohn–Sham NACs



From the interacting to the non-interacting KS representation

$$\psi^{mol.}(R, r) = \underline{\chi}^T(R) \underline{\Psi}^{elec.}(R, r)$$

**Adiabatic (interacting)
representation**

$$\left(\hat{T}_e + \hat{W}_{ee} + \hat{V}_{Ne}(R) \right) \Psi_\nu^{elec.}(R, r) = E_\nu^{elec.}(R) \Psi_\nu^{elec.}(R, r)$$

From the interacting to the non-interacting KS representation

$$\psi^{mol.}(R, r) = \underline{\chi}^T(R) \underline{\Psi}^{elec.}(R, r)$$

**Adiabatic (interacting)
representation**

$$= \underline{\chi}^T(R) \underline{\underline{\mathcal{U}}}(R) \underline{\underline{\mathcal{U}}}^\dagger(R) \underline{\Psi}^{elec.}(R, r)$$

*Unitary
transformation*

From the interacting to the non-interacting KS representation

$$\psi^{mol.}(R, r) = \underline{\chi}^T(R) \underline{\Psi}^{elec.}(R, r)$$

**Adiabatic (interacting)
representation**

$$= \underline{\chi}^T(R) \underline{\mathcal{U}}(R) \underline{\mathcal{U}}^\dagger(R) \underline{\Psi}^{elec.}(R, r) \longrightarrow \underline{\Phi}^{KS}(R, r)$$

From the interacting to the non-interacting KS representation

$$\psi^{mol.}(R, r) = \underline{\chi}^T(R) \underline{\Psi}^{elec.}(R, r)$$

**Adiabatic (interacting)
representation**

$$= \underline{\chi}^T(R) \underline{\mathcal{U}}(R) \underline{\mathcal{U}}^\dagger(R) \underline{\Psi}^{elec.}(R, r) \longrightarrow \underline{\Phi}^{KS}(R, r)$$

$$= \left(\underline{\mathcal{U}}^T(R) \underline{\chi}(R) \right)^T \underline{\Phi}^{KS}(R, r)$$

**Ensemble Kohn–Sham
representation**

Exact theory using Kohn–Sham non-adiabatic couplings

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) \left| \underline{\chi}(R) \right|^2 \right. \\ \left. -\frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \underline{\Lambda}_{\underline{\Phi}^{KS}}^{(j)}(R) \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \int dR \left| \underline{\chi}(R) \right|^2 \sum_{\nu \geq 0} \tilde{\xi}_\nu(R) E_\nu^{elec.}(R) \right\}$$

KS NACs
(no approximation
made!)



Exact theory using Kohn–Sham non-adiabatic couplings

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) \left| \underline{\chi}(R) \right|^2 \right. \\ \left. -\frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \underline{\Lambda}_{\underline{\Phi}_{KS}}^{(j)}(R) \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \int dR \left| \underline{\chi}(R) \right|^2 \sum_{\nu \geq 0} \tilde{\xi}_\nu(R) E_\nu^{elec.}(R) \right\}$$

Electronic ensemble energy
with modified weights

$$\tilde{\xi}_\nu(R) = \frac{\left| \sum_{\mu \geq 0} \mathcal{U}_{\nu\mu}^*(R) \chi_\mu(R) \right|^2}{\left| \underline{\chi}(R) \right|^2}$$

Now

$$\xi_\nu(R) \equiv \frac{|\chi_\nu(R)|^2}{|\underline{\chi}(R)|^2}$$

Before

Exact theory using Kohn–Sham non-adiabatic couplings

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) \left| \underline{\chi}(R) \right|^2 \right. \\ \left. -\frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \underline{\Lambda}_{\underline{\Phi}^{KS}}^{(j)}(R) \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \int dR \left| \underline{\chi}(R) \right|^2 \sum_{\nu \geq 0} \tilde{\xi}_\nu(R) E_\nu^{elec.}(R) \right\}$$

Unitary transformation

$$\underline{\Psi}^{elec.}(R, r) = \underline{\mathcal{U}}(R) \underline{\Phi}^{KS}(R, r)$$



$$\tilde{\xi}_\nu(R) = \frac{\left| \sum_{\mu \geq 0} \mathcal{U}_{\nu\mu}^*(R) \chi_\mu(R) \right|^2}{\left| \underline{\chi}(R) \right|^2}$$

Now

$$\xi_\nu(R) \equiv \frac{|\chi_\nu(R)|^2}{|\underline{\chi}(R)|^2}$$

Before

Exact theory using Kohn–Sham non-adiabatic couplings

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) \left| \underline{\chi}(R) \right|^2 \right. \\ \left. -\frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \underline{\Lambda}_{\underline{\Phi}^{KS}}^{(j)}(R) \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \int dR \left| \underline{\chi}(R) \right|^2 \sum_{\nu \geq 0} \tilde{\xi}_\nu(R) E_\nu^{elec.}(R) \right\}$$

Unitary transformation

$$\underline{\Psi}^{elec.}(R, r) = \underline{\mathcal{U}}(R) \underline{\Phi}^{KS}(R, r)$$

Ensemble energy

to be described with a KS system

Ensemble weights

$$\tilde{\xi}_\nu(R) = \frac{\left| \sum_{\mu \geq 0} \mathcal{U}_{\nu\mu}^*(R) \chi_\mu(R) \right|^2}{\left| \underline{\chi}(R) \right|^2}$$

Exact theory using Kohn–Sham non-adiabatic couplings

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) \left| \underline{\chi}(R) \right|^2 \right. \\ \left. -\frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \underline{\Lambda}_{\underline{\Phi}^{KS}}^{(j)}(R) \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \int dR \left| \underline{\chi}(R) \right|^2 \sum_{\nu \geq 0} \tilde{\xi}_\nu(R) E_\nu^{elec.}(R) \right\}$$

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$$\tilde{\xi}_\nu(R) = \frac{\left| \sum_{\mu \geq 0} \mathcal{U}_{\nu\mu}^*(R) \chi_\mu(R) \right|^2}{\left| \underline{\chi}(R) \right|^2} =$$

???



Exact theory using Kohn–Sham non-adiabatic couplings

$$\sum_{\nu} \tilde{\xi}_{\nu}(R) n_{\Psi_{\nu}^{elec.}}(R, \mathbf{r}) = n_{\underline{\xi}^{(R)}}(R, \mathbf{r}) = \sum_{\nu} \tilde{\xi}_{\nu}(R) n_{\Phi_{\nu}^{KS}}(R, \mathbf{r})$$

Ensemble density constraint

Unitary transformation

$$\underline{\Psi}^{elec.}(R, r) = \underline{\mathcal{U}}(R) \underline{\Phi}^{KS}(R, r)$$

Ensemble weights

$$\tilde{\xi}_{\nu}(R) = \frac{\left| \sum_{\mu \geq 0} \mathcal{U}_{\nu\mu}^*(R) \chi_{\mu}(R) \right|^2}{\left| \underline{\chi}(R) \right|^2} = ???$$

Exact theory using Kohn–Sham non-adiabatic couplings

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$$\underline{\Psi}^{elec.}(R, r) = \underline{\mathcal{U}}(R) \underline{\Phi}^{KS}(R, r)$$

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Functionals of the ensemble density!

Exact theory using Kohn–Sham non-adiabatic couplings

$$E = \min_{\underline{\chi}} \left\{ -\frac{1}{2M} \int dR \underline{\chi}^\dagger(R) \frac{\partial^2 \underline{\chi}(R)}{\partial R^2} + \int dR V_{NN}(R) \left| \underline{\chi}(R) \right|^2 \right. \\ \left. -\frac{1}{M} \sum_{j=1}^2 \frac{1}{j} \int dR \underline{\chi}^\dagger(R) \underline{\underline{\Lambda}}_{\underline{\Phi}^{KS}}^{(j)}(R) \frac{\partial^{2-j} \underline{\chi}(R)}{\partial R^{2-j}} \right. \\ \left. + \int dR \left| \underline{\chi}(R) \right|^2 \sum_{\nu \geq 0} \tilde{\xi}_\nu(R) E_\nu^{elec.}(R) \right\}$$

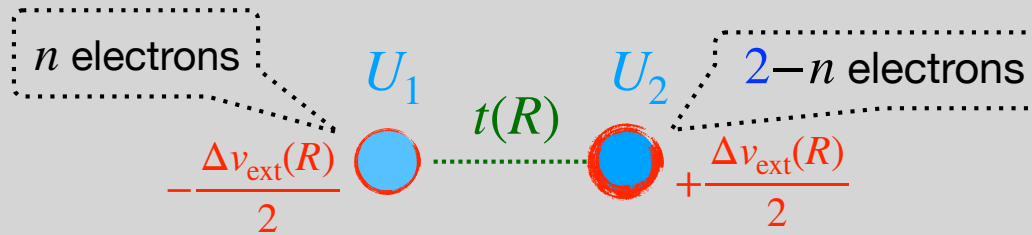
Unitary transformation

$$\underline{\Psi}^{elec.}(R, r) = \underline{\underline{\mathcal{U}}}(R) \underline{\Phi}^{KS}(R, r)$$

Self-consistent evaluation of the ensemble weights $\underline{\tilde{\xi}}(R)$

$$\underline{\tilde{\xi}}(R) = \underline{\tilde{\xi}} \left[\underline{\underline{\mathcal{U}}} \left[n_{\underline{\tilde{\xi}}(R)} \right], \underline{\chi}(R) \right]$$

Implementation of the theory (work in progress)

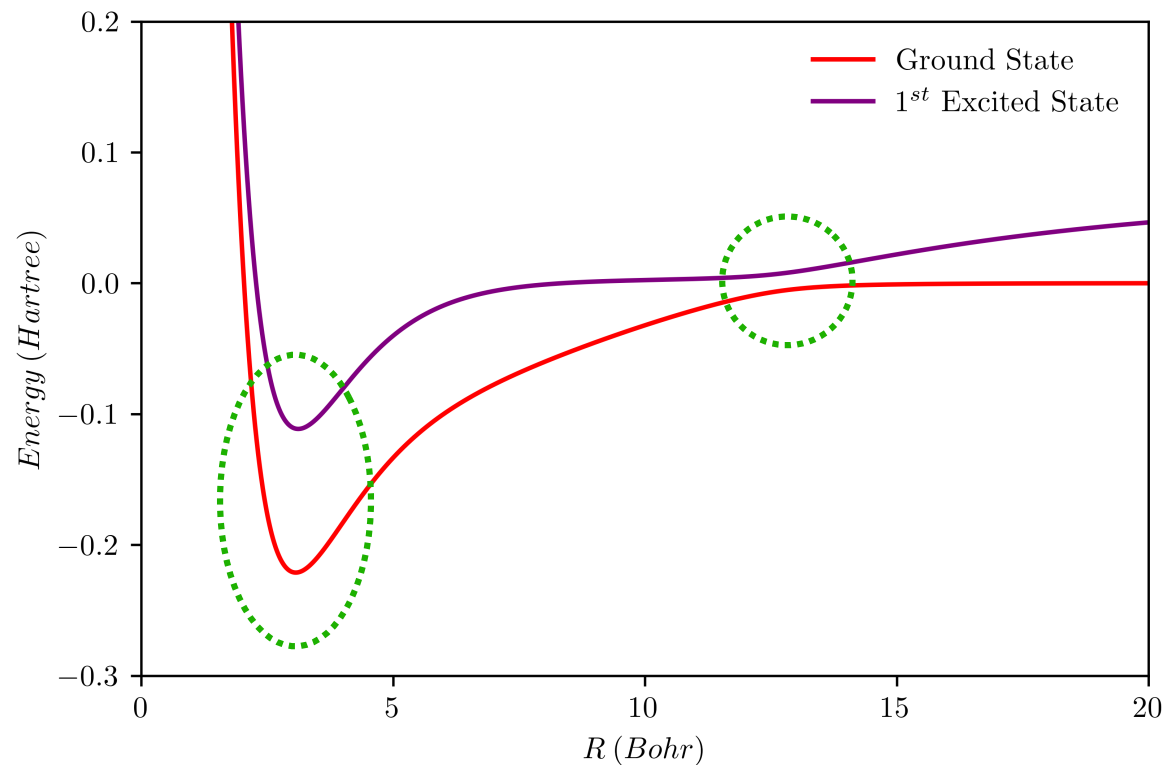


Geometry-dependent
two-electron **Hubbard dimer model**

THE JOURNAL OF CHEMICAL PHYSICS **148**, 084110 (2018)

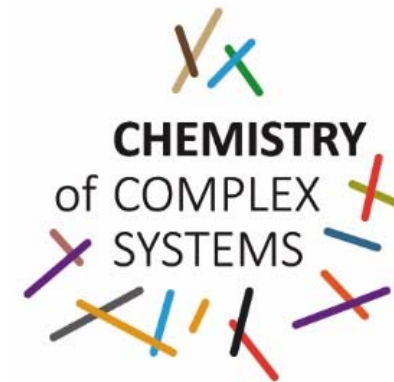
Density functional theory of electron transfer beyond the Born-Oppenheimer approximation: Case study of LiF

Chen Li,^{1,a)} Ryan Requist,¹ and E. K. U. Gross^{1,2}



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“Lab of Excellence” project:
LabEx CSC (ANR-10-LABX-0026-CSC)



« Initiative d'Excellence » (IdEx) grant
University of Strasbourg



CoLab ANR project



Unified N -centered ensemble DFT of charged and neutral excited states

$$n(\mathbf{r}) \equiv \left(1 - \sum_{\nu>0}^{\text{excited states}} \frac{N_\nu}{N} \xi_\nu \right) n_0^N(\mathbf{r}) + \sum_{\nu>0}^{\text{excited states}} \xi_\nu n_\nu(\mathbf{r})$$

Ensemble enlarged to **charged** excitations

B. Senjean and E. Fromager, *Phys. Rev. A* **98**, 022513 (2018).


B. Senjean and E. Fromager, *Int. J. Quantum Chem.* 2020; 120:e26190

F. Cernatic, B. Senjean, V. Robert, and E. Fromager, *Top Curr Chem (Z)* **380**, 4 (2022).

C. Marut, F. Cernatic, B. Senjean, P.-F. Loos, and E. Fromager, to be submitted (2023).

Unified N -centered ensemble DFT of charged and neutral excited states

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$$\int d\mathbf{r} n(\mathbf{r}) = N$$

B. Senjean and E. Fromager, *Phys. Rev. A* **98**, 022513 (2018).

B. Senjean and E. Fromager, *Int. J. Quantum Chem.* 2020; 120:e26190

F. Cernatic, B. Senjean, V. Robert, and E. Fromager, *Top Curr Chem (Z)* **380**, 4 (2022).

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Unified N -centered ensemble DFT of charged and neutral excited states

$$n(\mathbf{r}) \equiv \left(1 - \sum_{\nu>0}^{\text{excited states}} \left(\frac{N_\nu}{N} \right)^{\xi_\nu} \right) n_0^N(\mathbf{r}) + \sum_{\nu>0}^{\text{excited states}} \xi_\nu n_\nu(\mathbf{r})$$



$$E\{\xi_\nu\}[n] \equiv \left(1 - \sum_{\nu>0}^{\text{excited states}} \left(\frac{N_\nu}{N} \right)^{\xi_\nu} \right) E_0^N + \sum_{\nu>0}^{\text{excited states}} \xi_\nu E_\nu$$

B. Senjean and E. Fromager, *Phys. Rev. A* **98**, 022513 (2018).

B. Senjean and E. Fromager, *Int. J. Quantum Chem.* 2020; 120:e26190

F. Cernatic, B. Senjean, V. Robert, and E. Fromager, *Top Curr Chem (Z)* **380**, 4 (2022).

C. Marut, F. Cernatic, B. Senjean, P.-F. Loos, and E. Fromager, to be submitted (2023).

$$E_{\mu}^N - E_0^N = \left(E_{\mu}^N - E_0^{N-1} \right) + \left(E_0^{N-1} - E_0^N \right)$$

$$E_0^{N-1} - E_0^N = -\varepsilon_N^{\xi_{\mu}=0}$$

$$+ \frac{\left(\int d\mathbf{r} v_{\text{Hxc}}^{\xi_{\mu}=0}(\mathbf{r}) n_0^N(\mathbf{r}) - E_{\text{Hxc}}[n_0^N] \right)}{N} + \frac{\partial E_{\text{xc}}^{\{\xi_{\nu}\}}[n_0^N]}{\partial \xi_{-}} \Big|_{\{\xi_{\nu}\}=0}$$

$$\xi_{\mu} = 0$$

Koopmans theorems

if $\equiv 0$ (arbitrary choice)

$$\xi_{\mu} \rightarrow 0^+$$

$$E_{\mu}^N - E_0^{N-1} = \varepsilon_{N+\mu}^{\xi_{\mu} \rightarrow 0^+}$$

$$- \frac{\left(\int d\mathbf{r} v_{\text{Hxc}}^{\xi_{\mu} \rightarrow 0^+}(\mathbf{r}) n_0^N(\mathbf{r}) - E_{\text{Hxc}}[n_0^N] \right)}{N} - \frac{\partial E_{\text{xc}}^{\{\xi_{\nu}\}}[n_0^N]}{\partial \xi_{-}} \Big|_{\{\xi_{\nu}\}=0} + \frac{\partial E_{\text{xc}}^{\{\xi_{\nu}\}}[n_0^N]}{\partial \xi_{\mu}} \Big|_{\{\xi_{\nu}\}=0}$$

$$E_{\mu}^N - E_0^N = \varepsilon_{N+\mu}^{\xi_{\mu} \rightarrow 0^+} - \varepsilon_N^{\xi_{\mu}=0}$$

**Exact neutral
excitation energy**

M. Levy, *Phys. Rev. A* **52**, R4313 (1995).

T. Gould, Z. Hashimi, L. Kronik, and S. G. Dale, *J. Phys. Chem. Lett.* **13**, 2452 (2022).

C. Marut, F. Cernatic, B. Senjean, P.-F. Loos, and E. Fromager, to be submitted (2023).

$$E_0^{N-1} - E_0^N = -\varepsilon_N^{\xi_\mu=0} + \frac{\left(\int d\mathbf{r} v_{\text{Hxc}}^{\xi_\mu=0}(\mathbf{r}) n_0^N(\mathbf{r}) - E_{\text{Hxc}}[n_0^N] \right)}{N} + \left. \frac{\partial E_{\text{xc}}^{\{\xi_\nu\}}[n_0^N]}{\partial \xi_-} \right|_{\{\xi_\nu\}=0}$$

$\xi_\mu = 0$

Koopmans theorems

if $\xi_\mu = 0$ (arbitrary choice)

$$E_\mu^N - E_0^{N-1} \underset{\mu > 0}{=} \varepsilon_{N+\mu}^{\xi_\mu \rightarrow 0^+} - \frac{\left(\int d\mathbf{r} v_{\text{Hxc}}^{\xi_\mu \rightarrow 0^+}(\mathbf{r}) n_0^N(\mathbf{r}) - E_{\text{Hxc}}[n_0^N] \right)}{N} - \left. \frac{\partial E_{\text{xc}}^{\{\xi_\nu\}}[n_0^N]}{\partial \xi_-} \right|_{\{\xi_\nu\}=0} + \left. \frac{\partial E_{\text{xc}}^{\{\xi_\nu\}}[n_0^N]}{\partial \xi_\mu} \right|_{\{\xi_\nu\}=0}$$

$\xi_\mu \rightarrow 0^+$

$$\int \frac{d\mathbf{r}}{N} \left(v_{\text{xc}}^{\xi_\mu \rightarrow 0^+}(\mathbf{r}) - v_{\text{xc}}^{\xi_\mu=0}(\mathbf{r}) \right) n_0^N(\mathbf{r}) = \left. \frac{\partial E_{\text{xc}}^{\{\xi_\nu\}}[n_0^N]}{\partial \xi_\mu} \right|_{\{\xi_\nu\}=0}$$

Derivative discontinuity!

Extraction of individual energies from the ensemble energy

$$E\{\xi_\nu\} = \left(1 - \sum_{\nu>0} \xi_\nu\right) E_0 + \sum_{\nu>0} \xi_\nu E_\nu$$



Auxiliary quantity (not an observable) ...

Extraction of individual energies from the ensemble energy

$$E\{\xi_\nu\} = \left(1 - \sum_{\nu>0} \xi_\nu\right) E_0 + \sum_{\nu>0} \xi_\nu E_\nu$$



Auxiliary quantity (not an observable) ...

$$E\{\xi_\nu\} = \sum_{\mu \geq 0} \xi_\mu \langle \hat{T} + \hat{V}_{\text{ext}} \rangle_{\Phi_\mu^{\{\xi_\nu\}}} + E_{\text{Hxc}}^{\{\xi_\nu\}} \left[n^{\{\xi_\nu\}} \right]$$

*We cannot “read”
the individual energies!*

Extraction of individual energies from the ensemble energy

$$E\{\xi_\nu\} = \left(1 - \sum_{\nu>0} \xi_\nu\right) E_0 + \sum_{\nu>0} \xi_\nu E_\nu$$



Auxiliary quantity (not an observable) ...

*... that **varies linearly** with the ensemble weights!*

Extraction of individual energies from the ensemble energy

$$E\{\xi_\nu\} = \left(1 - \sum_{\nu>0} \xi_\nu\right) E_0 + \sum_{\nu>0} \xi_\nu E_\nu$$



Auxiliary quantity (not an observable) ...

*... that **varies linearly** with the ensemble weights!*

$$E_\mu = E_0 + E_\mu - E_0$$

Extraction of individual energies from the ensemble energy

$$E\{\xi_\nu\} = \left(1 - \sum_{\nu>0} \xi_\nu\right) E_0 + \sum_{\nu>0} \xi_\nu E_\nu$$



Auxiliary quantity (not an observable) ...

... that varies linearly with the ensemble weights!

$$E_\mu = E_0 + E_\mu - E_0 = E_0 + \sum_{\nu>0} \delta_{\nu\mu} \frac{\partial E\{\xi_\nu\}}{\partial \xi_\nu}$$

Extraction of individual energies from the ensemble energy

$$E\{\xi_\nu\} = \left(1 - \sum_{\nu>0} \xi_\nu\right) E_0 + \sum_{\nu>0} \xi_\nu E_\nu$$



Auxiliary quantity (not an observable) ...

... that varies linearly with the ensemble weights!

$$E_\mu = E_0 + E_\mu - E_0 = E_0 + \sum_{\nu>0} \delta_{\nu\mu} \frac{\partial E\{\xi_\nu\}}{\partial \xi_\nu}$$

$$E\{\xi_\nu\} = 0$$

Extraction of individual energies from the ensemble energy

$$E\{\xi_\nu\} = \left(1 - \sum_{\nu>0} \xi_\nu\right) E_0 + \sum_{\nu>0} \xi_\nu E_\nu$$



Auxiliary quantity (not an observable) ...

... that varies linearly with the ensemble weights!

$$E_\mu = E_0 + E_\mu - E_0 = E_0 + \sum_{\nu>0} \delta_{\nu\mu} \frac{\partial E\{\xi_\nu\}}{\partial \xi_\nu} = E\{\xi_\nu\} + \sum_{\nu>0} \left(-\xi_\nu + \delta_{\nu\mu}\right) \frac{\partial E\{\xi_\nu\}}{\partial \xi_\nu}$$

(Note: In the original image, E_0 , the sum term, and the derivative term are enclosed in purple rounded rectangles. A purple arrow points from the E_0 box to a grey box containing $E\{\xi_\nu\}=0$. Another purple arrow points from this grey box to the sum term box. A third purple arrow points from the grey box to the derivative term box.)

Extraction of individual energies from the ensemble energy

$$E\{\xi_\nu\} = \left(1 - \sum_{\nu>0} \xi_\nu\right) E_0 + \sum_{\nu>0} \xi_\nu E_\nu$$



Auxiliary quantity (not an observable) ...

*... that **varies linearly** with the ensemble weights!*

$$E_\mu = E\{\xi_\nu\} + \sum_{\nu>0} \left(\delta_{\nu\mu} - \xi_\nu\right) \frac{\partial E\{\xi_\nu\}}{\partial \xi_\nu}$$

Extraction of individual energies from the ensemble energy

$$E\{\xi_\nu\} = \left(1 - \sum_{\nu>0} \xi_\nu\right) E_0 + \sum_{\nu>0} \xi_\nu E_\nu$$



Auxiliary quantity (not an observable) ...

... that varies linearly with the ensemble weights!

$$E_\mu = \left[1 + \sum_{\nu>0} \left(\delta_{\nu\mu} - \xi_\nu\right) \frac{\partial}{\partial \xi_\nu}\right] E\{\xi_\nu\}$$

**Ensemble Kohn-Sham
decomposition**

Extraction of individual energies from the ensemble energy

$$E\{\xi_\nu\} = \left(1 - \sum_{\nu>0} \xi_\nu\right) E_0 + \sum_{\nu>0} \xi_\nu E_\nu$$



Auxiliary quantity (not an observable) ...

... that varies linearly with the ensemble weights!

$$E_\mu = \left[1 + \sum_{\nu>0} \left(\delta_{\nu\mu} - \xi_\nu \right) \frac{\partial}{\partial \xi_\nu} \right] E\{\xi_\nu\}$$

Deduced from the (ensemble) Kohn-Sham orbital energies

Ensemble Kohn-Sham decomposition

Ordered-weight ensemble density functional theory

$$E\{\xi_\nu(R)\} = \sum_{\nu \geq 0} \xi_\nu(R) E_\nu^{elec.}(R)$$

What if the weights are **not**
in decreasing order?

Ordered-weight ensemble density functional theory

$$E\{\xi_\nu(R)\} = \sum_{\nu \geq 0} \xi_\nu(R) E_\nu^{elec.}(R)$$



$$E[\xi_\nu(R)] = \sum_{\nu \geq 0} [\xi_\nu(R)] E_\nu^{elec.}(R)$$

*Ordered-weight
Ensemble energy*

Ordered-weight ensemble density functional theory

$$E\{\xi_\nu(R)\} = \sum_{\nu \geq 0} \xi_\nu(R) E_\nu^{elec.}(R)$$



$$E[\xi_\nu(R)] = \sum_{\nu \geq 0} [\xi_\nu(R)] E_\nu^{elec.}(R)$$

*Ordered-weight
Ensemble energy*

$$n[\xi_\nu(R)](\mathbf{r}) = \sum_{\nu \geq 0} [\xi_\nu(R)] n_{\Psi_\nu^{elec.}(R)}(\mathbf{r})$$

*Ordered-weight
Ensemble density*

Ordered-weight ensemble density functional theory

$$E\{\xi_\nu(R)\} = \sum_{\nu \geq 0} \xi_\nu(R) E_\nu^{elec.}(R)$$



$$E[\xi_\nu(R)] = \sum_{\nu \geq 0} [\xi_\nu(R)] E_\nu^{elec.}(R)$$

$$n[\xi_\nu(R)](\mathbf{r}) = \sum_{\nu \geq 0} [\xi_\nu(R)] n_{\Psi_\nu^{elec.}(R)}(\mathbf{r})$$



$$E_\nu^{elec.}(R) \equiv E_\nu^{elec.} [n[\xi_\nu(R)]]$$

Ordered-weight ensemble density functional theory

